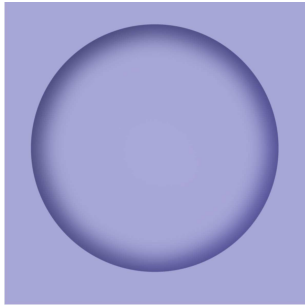


Macroclumping in winds of massive stars

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Clumping: micro versus macro

The "microclumping" approximation:
 (included in various codes: CMFGEN, PoWR, Munich codes ...)

Assumptions

- Clumps are *optically thin at all frequencies*
- For frequencies of large opacity, this implies unrealistically small scales of the clumps
- In the winds of O- and WR stars, high opacities occur especially
 - in the resonance lines
 - in the X-ray domain (K-shell absorption)

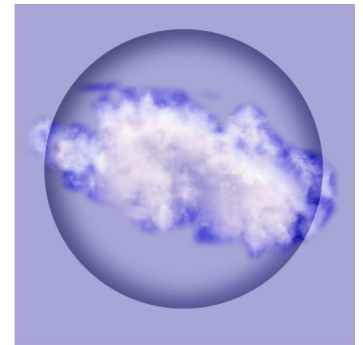
Photon mean free path

maximum size for opt. thin clumps
Example: ζ Puppis model
 at typical point in the wind (at $\frac{1}{2}v_\infty$)

C IV resonance line	$4 \cdot 10^{-4} R_*$
P V resonance line	$0.1 R_*$
He II 1 - 2 at 304 Å	$3 \cdot 10^{-6} R_*$
K-shell edges (X-ray)	$2 \cdot 10^{-4} R_*$

An approach to macroclumping (Oskinova, Hamann & Feldmeier 2007)

- A large number of clumps are statistically distributed
- At a given frequency and location, all clumps have the optical diameter τ_{clump}
- Macroclumping reduces the effective opacity by a factor $\frac{1 - e^{-\tau_{\text{clump}}}}{\tau_{\text{clump}}}$
- For $\tau_{\text{clump}} \gg 1$, this implies a drastic reduction by a factor of τ_{clump}^{-1} !
- The reduction of opacity is already significant for moderate $\tau_{\text{clump}} \approx 1$
- Specification of the clump size:
 - Clumps are conserved entities (\rightarrow eq. of continuity)
 - the average separation of clumps is, with a parameter L_0 , $L(r) = L_0 \left(r^2 \frac{v(r)}{v_\infty} \right)^{1/3}$
- Realistic choices for L_0 are of the order of $1 R_*$
 - from 1-D hydro (Feldmeier et al. 1997): radial shells launched on \lesssim dynamical timescale $\rightarrow L_0 \lesssim 1 R_*$
 - from *Line Profile Variability* in WR stars: $\gtrsim 10^4$ clumps launched per hour (Lépine & Moffatt 1999) $\rightarrow L_0 \approx 0.25 R_*$
 - from modeling X-ray line profiles (Oskinova et al. 2006): $L_0 \lesssim 1 R_*$



Macroclumping can mask higher mass-loss rates

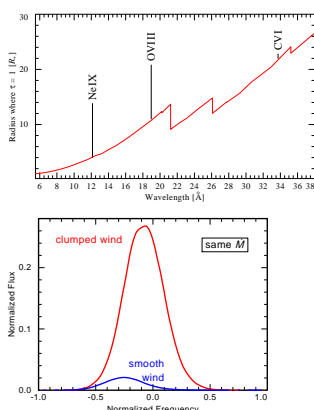
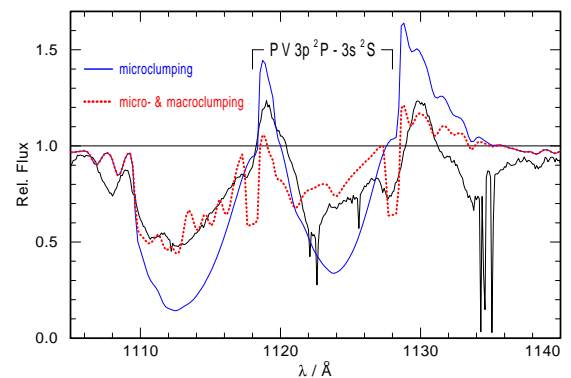
Implementation in the Potsdam Wolf-Rayet (PoWR) code

- Macroclumping only accounted in the *Formal Integral*
- *Effective opacity*, reduced due to macroclumping effect
- Non-LTE source function as from microclumping model, i.e. feedback of macroclumping effects on population numbers neglected

Macroclumping can solve the mass-loss rate discrepancy

- The P V resonance line gives >10 times smaller O-star mass-loss rates than H α and radio free-free emission! (Fullerton, Massa & Prinja 2006)
- Accounting only for *microclumping* \rightarrow incredibly small filling factors to reconcile the p^2 diagnostics (H α line, radio free-free emission) with such small mass loss
- Accounting for *macroclumping* effect \rightarrow weaker P V resonance line, reconciling it with the p^2 diagnostics without a reduction of the mass-loss rate

Figure: Observation of ζ Pup and two synthetic profiles from the same model, but with and without macroclumping effect, respectively



Macroclumping explains how X-rays from massive-star winds escape

- In soft X-rays, stellar winds should be opaque due to K-shell absorption (*left Figure, upper panel*: radius of optical depth unity for a ζ Pup model)
- From the X-ray line widths ($\frac{1}{2}v_\infty$, typically) and from hydrodynamical models, X-rays are expected to be produced in the lower part of the wind
- The reduction of the effective opacity by *macroclumping* can explain how a significant fraction of the X-rays can emerge (*left Figure, lower panel*)
- When clumps are optically thick, the effective opacity becomes *grey*; this can explain the observed similarity between different X-ray line profiles
- Observed *symmetric* X-ray line profiles can be reproduced best with macroclumps that are flat like pancakes ("broken-shell model", *right Figure*) Oskinova et al. 2004

